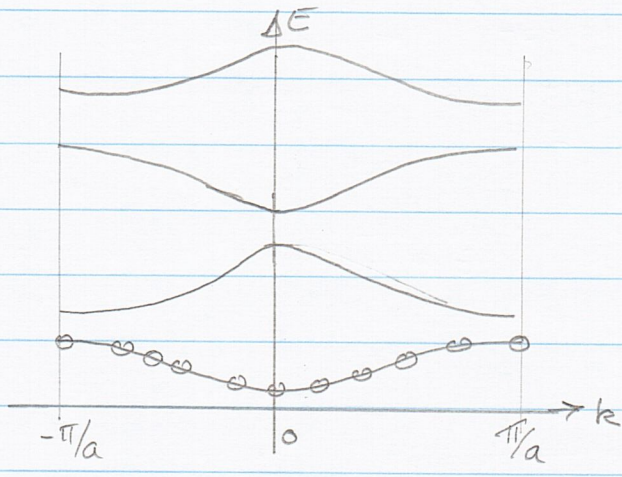


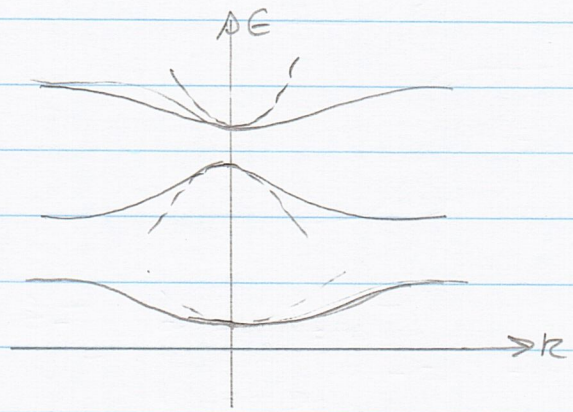
Pre-MT summary EEE 480 2011

LECT. 2.



Energy-band structure
 Reduced-zone plot
 States ($2N/\text{band}$)
 Bandgap

LECT. 3



Parabolic band approximation
 $E_c = E_{c0} + \frac{\hbar^2 k^2}{2m_e^*}$ for electrons

$$m_e^* \propto \left[\frac{d^2 E}{dk^2} \right]^{-1}$$

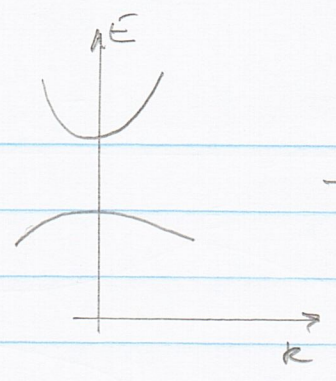
$$F_{ext} = m_e^* \frac{dv}{dt}$$

Factors affected by m_e^* : $\mu \propto 1/m_e^* \rightarrow$ more easily accelerated
 $N_c \propto m_e^* \rightarrow$ more states/ ΔE

Factors affected by E_g : $n_i \rightarrow$ more E needed to break bond
 $J_0 \downarrow (\propto n_i^2)$
 $V_{oc} \uparrow (\propto \ln J_{ph}/J_0)$
 $J_{ph} \downarrow$ (more transparent)

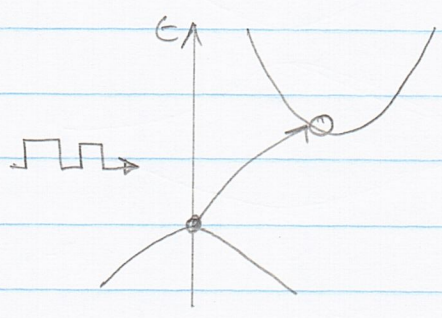
LECT. 4

Band diagram

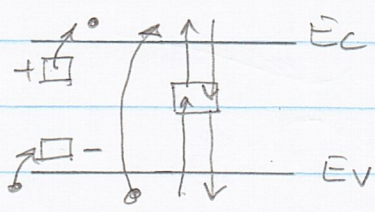


$$E_c = E_{c0} + U_M(x)$$

$$E_v$$



Indirect bandgap material
phonons



Intrinsic Extrinsic
Ac Donors, Acceptors
R-G centres

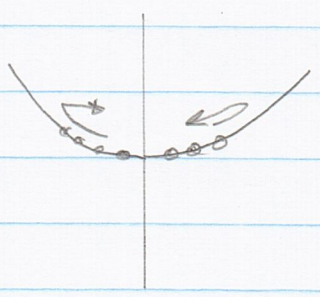
Thermal Equilib $U = R_{th} - G_{th} = 0$

$$R_{PG} = \bar{A}n \quad (\text{p-type material})$$

$$U_{PG} = A(n_0 + \Delta n) - An_0 = A\Delta n$$

$$\stackrel{\Delta}{=} \Delta n \quad (\text{p-type}).$$

$$\bar{n}_{e,eq} = f(N_A, N_D) \begin{matrix} (3.21) \\ (3.22) \end{matrix}$$



Scattering - randomizing of velocities
- keeps carriers near band extrema.

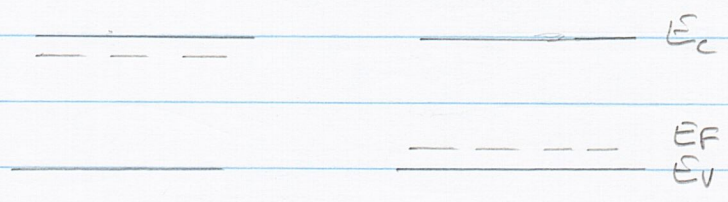
LECT 5



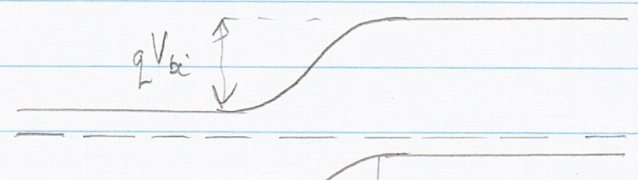
Fermi-Dirac distⁿ for state filling

$$f \cdot (E) = \frac{1}{1 + \exp^{E - E_F / k_B T}}$$

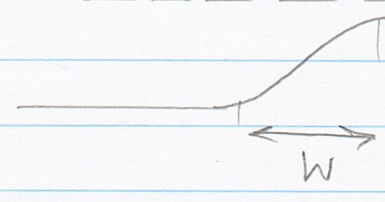
LECT 5 (cont)



Fermi level

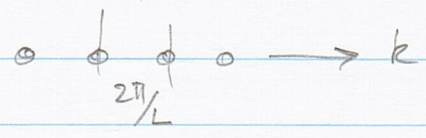


Band diagram (Equilib)



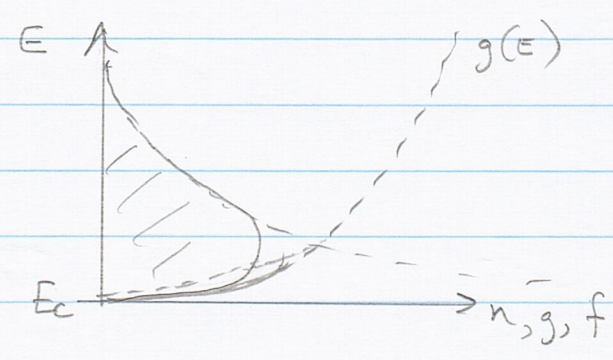
Built-in voltage
Depletion region

LECT 6



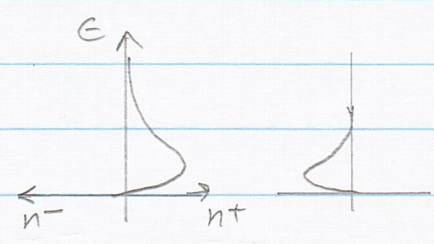
Density of states $g(E)$

$$n_0 = \int_{E_c} g(E) f(E) dE \rightarrow N_c \exp \frac{E_F - E_c}{k_B T} \text{ for MB dist}^n$$

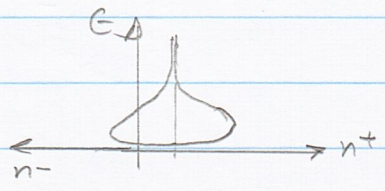


Hemi-Maxwellian

$$\langle v_x^2 \rangle = 2v_T^2 \propto \frac{1}{m_e^*}$$



Diffusion: $J_e = q D_e \frac{dn}{dx}$



Displaced Maxwellian

$$J_e = -q n v_d = +q n \mu E$$

LECT. 7

$$J_e = q n \mu E + q D \frac{dn}{dx}$$

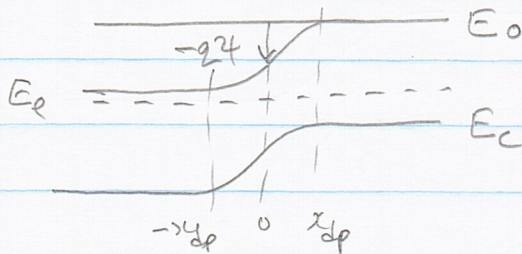
DDE

$$D_e = \frac{kT}{q} \mu$$

Einstein Relation

$$\mu = f(N_A, N_D)$$

(5.30, 5.31)



Valence levels

$$-q\phi$$

$$n_0 \propto \exp\left(-\frac{q\phi}{kT}\right)$$

$$-\frac{d^2\phi}{dx^2} = \frac{\rho}{\epsilon} = \frac{q}{\epsilon} (-n + p - N_A + N_D) \quad n, p \text{ f } (\phi)$$

$$\left. \begin{aligned} n(x), p(x) &\ll N_D \quad -x_p \leq x \leq 0 \\ &\ll N_A \quad 0 \leq x \leq x_p \end{aligned} \right\} \text{Depletion Approximation}$$

e.g. $-\frac{d^2\phi}{dx^2} = -\frac{qN_D}{\epsilon} \quad -x_n \leq x \leq 0$

BC's $-\frac{d\phi}{dx} = 0$ @ $x = -x_n$
 $-\frac{d\phi}{dx} = 0$ @ $x = x_p$
 $\phi = 0$ @ $x = x_p$

$$\left. \begin{aligned} & \text{Find } \phi(-x_n) - \phi(x_p) \\ & = V_J \\ & = V_{bi} \quad (V_a = 0) \\ & = V_{bi} - V_a \quad (\text{biased}) \end{aligned} \right\}$$

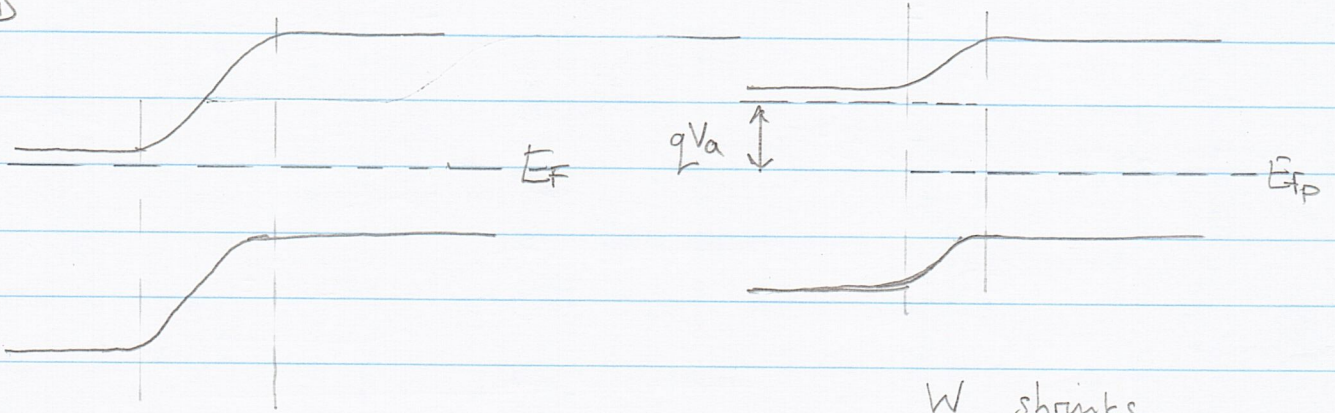
$$\phi(x) = \frac{qN_D}{2\epsilon} \left(\frac{x^2}{2} - x_n x \right) \quad -x_n \leq x \leq 0 \quad \phi = f(x)$$

$$\begin{aligned} qN_D x_n &= qN_A x_p \\ x_n + x_p &= W \end{aligned} \quad \text{charge neutrality}$$

$$W = \sqrt{\frac{2\epsilon}{q} V_J \left(\frac{1}{N_A} + \frac{1}{N_D} \right)}$$

LECT. 8

FORWARD
BIAS



W shrinks

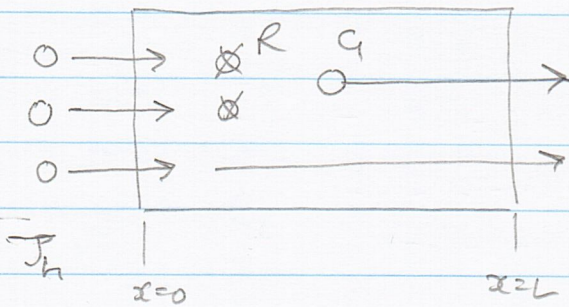
$$E_{Fn}(0) - E_{Fp}(L) = qV_a$$

$$qV_a = V_J = (V_{bi} + V_a)$$

$E_{Fn}, E_{Fp} \approx \text{flat across } W$

$$n = N_c \exp\left(\frac{E_{Fn} - E_c}{k_B T}\right) \equiv n_i \exp\left(\frac{E_{Fn} - E_{Fi}}{k_B T}\right)$$

$$E_{Fn} \triangleq \frac{k_B T \ln n}{N_c} + E_c$$

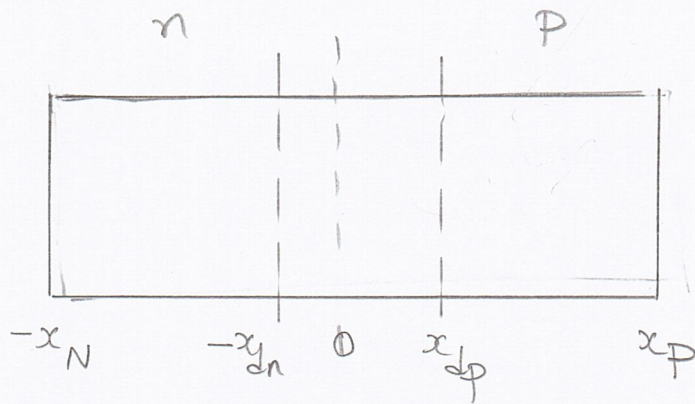


Continuity Equation (for holes)

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot J_h - R + G$$

$$\underbrace{-R + G}_{-(R-G)} \triangleq -U_h$$

Recall $U_h = \frac{\Delta p}{\tau_h} = \frac{p - p_{on}}{\tau_h}$



Sec. 5.3

The master equation set

8

This is our version
of (5.24)

$$\begin{aligned}
 -\nabla^2 \psi &= \frac{q}{\epsilon} [p - n + N_D - N_A] \\
 J_e &= -qn\mu_e \nabla \psi + qD_e \nabla n \\
 J_h &= -qp\mu_h \nabla \psi - qD_h \nabla p \\
 \frac{\partial n}{\partial t} &= \frac{1}{q} \nabla \cdot J_e - \frac{n - n_0}{\tau_e} + G_{op} \\
 \frac{\partial p}{\partial t} &= -\frac{1}{q} \nabla \cdot J_h - \frac{p - p_0}{\tau_h} + G_{op}
 \end{aligned}$$

J 's are $f(n,p)$, so, essentially, we have

3 equations in 3 unknowns ;

solve numerically.

OR simplify & solve
analytically

BCs: $x = -x_N$: $p = p_{on}$ OHMIC ; $J_h \frac{dp}{dx} = S_F (p(-x_N) - p_{on})$

$x = x_p$: $n = n_{op}$ " ; $-J_e \frac{dn}{dx} = S_B (n(x_p) - n_{op})$

$x = x_{dp}$: $n = n_{op}$ SOLAR CELL PHOTOCURRENT

$x = x_{dp}$: $n = n_{op} \exp^{-qV_a/k_B T}$ DARK DIODE

$x = -x_{dn}$: $p = p_{on} \exp^{-qV_a/k_B T}$ DARK DIODE

In p-n junction, e.g.,

$$J_e = q D_e \frac{dn}{dx}$$

$$0 = \frac{1}{q} \frac{dJ_e}{dx} - \frac{(n - n_{op})}{\tau_e}$$

$$\rightarrow 0 = \frac{d^2 n}{dx^2} - \frac{(n - n_{op})}{L_e^2} \quad L_e = \sqrt{D_e \tau_e}$$

General solution: $n(x) - n_{op} = A e^{x/L_e} + B e^{-x/L_e}$

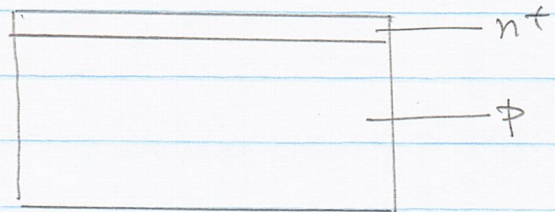
LECT 9

SOLAR CELL: Add q_{op} to continuity equations

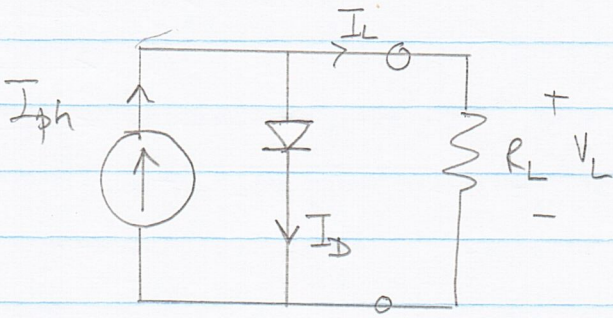
$$q_{op} = \alpha \Phi_0 e^{-\alpha x}$$

α = absorption coeff

Φ = photon flux = $f(\lambda)$

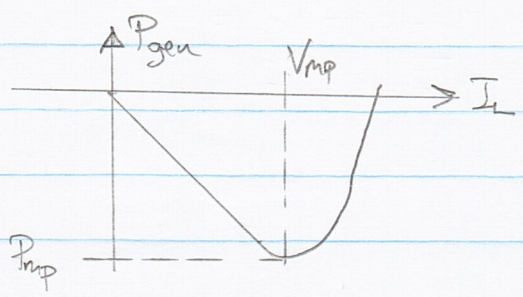
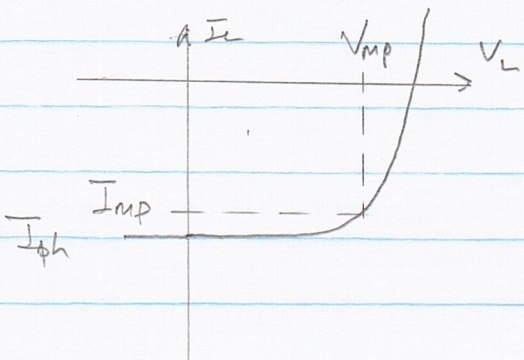


LECT 10



$$I_L = I_{ph} - I_0 \left(\exp \frac{qV}{k_B T} - 1 \right)$$

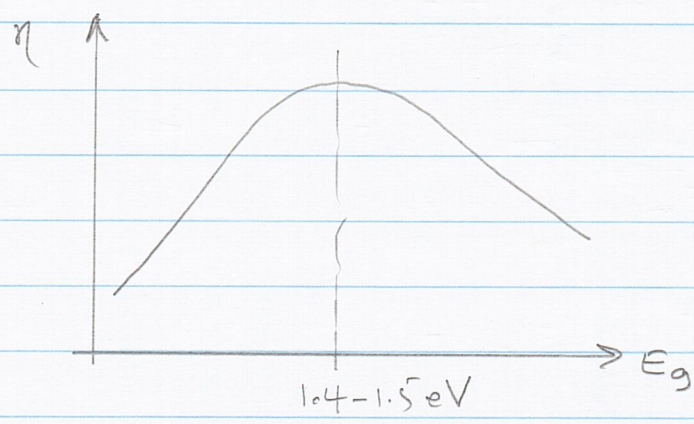
$$V_{oc} = \frac{k_B T}{q} \ln \frac{I_{ph} + I_0}{I_0}$$



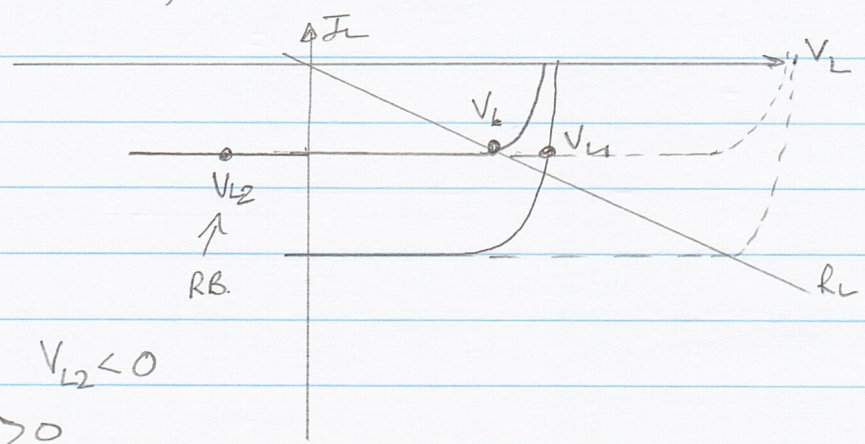
$$P_{mp} = FF J_{sc} V_{oc}$$

LECT 11

$$\eta = \frac{FF J_{sc} V_{oc}}{P_{in}}$$



TWO CELLS IN SERIES, ONE SHADED



$$I_L < 0 \quad V_{L2} < 0$$

$$P_{L2} > 0$$